

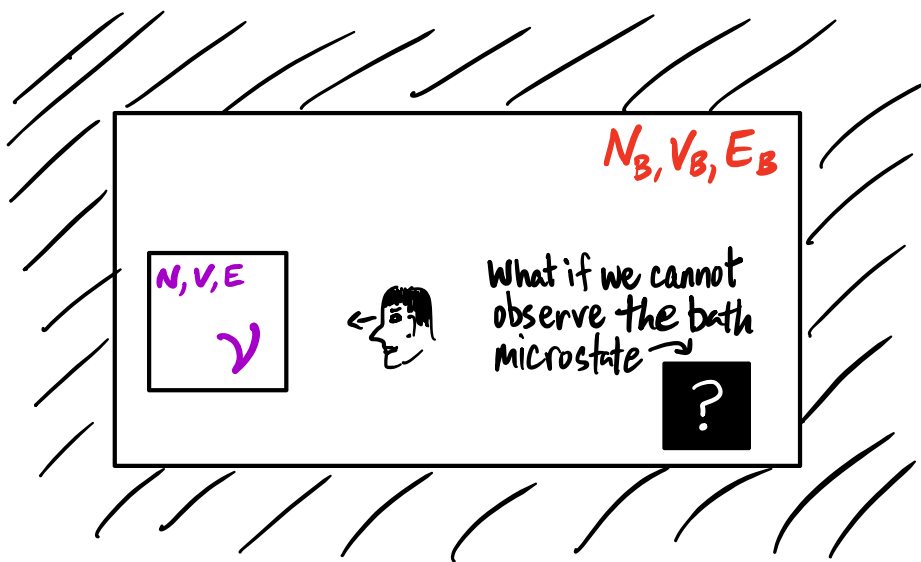
Lecture 3

Recall from last lecture...

System + (Thermal) Bath

Microstate of the full system + bath is defined by

$$\nu_T = \{ \nu, \nu_B \}$$



Variables with no flow:

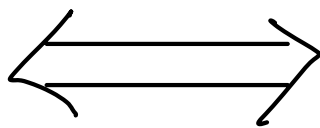
$$N_T = N + N_B$$
$$V_T = V + V_B$$

Variable with flow:

$$E_T = E(\nu) + E_B(\nu_B)$$

System microstate Bath microstate

Probability of
system microstate
 ν



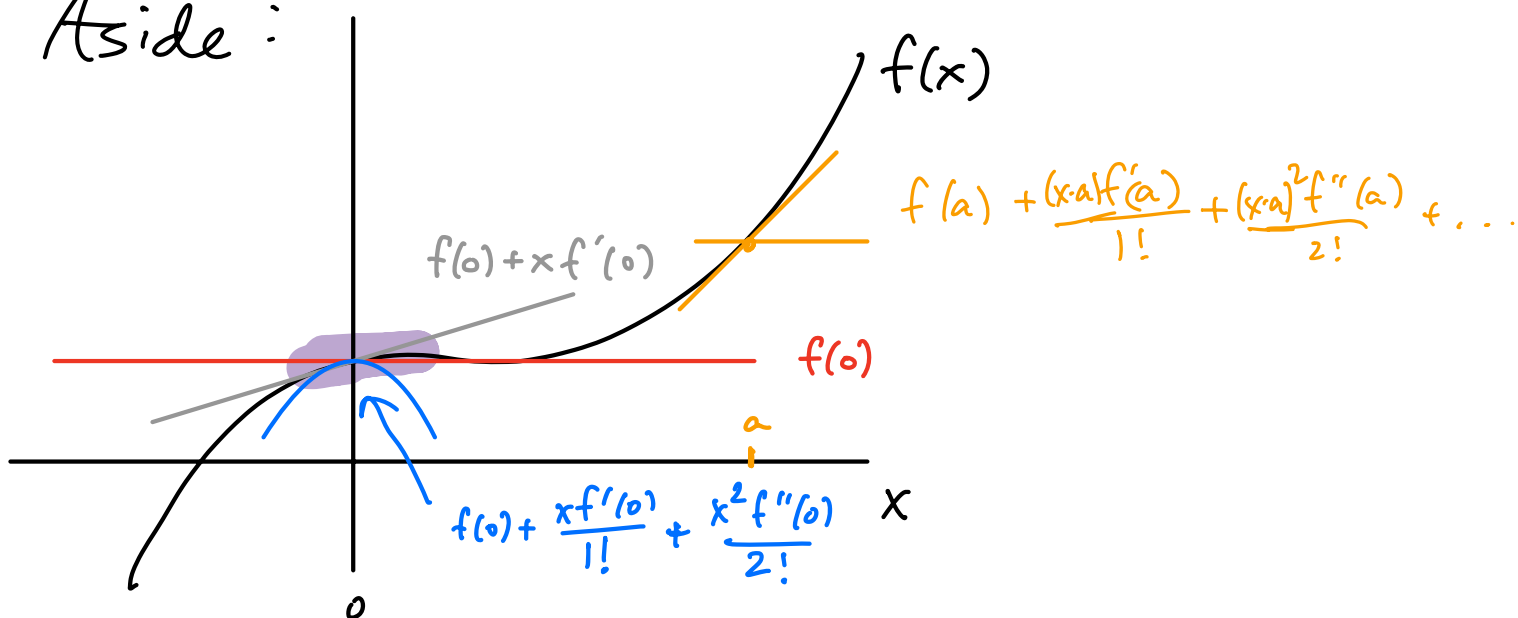
How many bath
microstates have
energy $E_T - E(\nu)$

$$P(\nu) \propto \Omega_B(E_T - E(\nu)) \quad (\text{Marginalization})$$

What is $\Omega_B(E_T - E)$ in the limit of small E ?

Taylor Expand!

Aside:



(Let's actually Taylor expand $\ln \Omega_B(E_T - E)$ not $\Omega_B(E_T - E)$)

$$\ln \Omega_B(E_T - E) \approx \ln \Omega_B(E_T) + E \left(\frac{\partial \ln \Omega_B}{\partial E} \right) \bigg|_{E=0} + \dots$$

How do the # of bath microstates change with a change in system energy E ?
 ^
 (on a log scale)

Notice that $E_T - E = E_B \Rightarrow$

$$\left(\frac{\partial \ln \Omega_B}{\partial E} \right) \bigg|_{E=0} = \left(\frac{\partial \ln \Omega_B}{\partial E_B} \right) \underbrace{\left(\frac{\partial E_B}{\partial E} \right)}_{\Downarrow -1} \quad \text{Chain Rule}$$
$$= - \left(\frac{\partial \ln \Omega_B}{\partial E_B} \right)_{N_B, V_B}$$

The BIG INSIGHT...

$\left(\frac{\partial \ln \Omega_B}{\partial E_B} \right)_{N_B, V_B}$ has nothing to do with the system.
It is a property of the bath!

Let's just define (for now) that bath property as

$$\beta = \left(\frac{\partial \ln \Omega_B}{\partial E_B} \right)_{N_B, V_B}$$

The Taylor expansion becomes

$$\ln \Omega_B(E_T - E) \approx \ln \Omega_B(E_T) + E \left(\frac{\partial \ln \Omega_B}{\partial E} \right) \Big|_{E=0} + \dots$$

$$\ln \Omega_B(E_T - E) \approx \ln \Omega_B(E_T) - \beta E + \dots ?$$

chain

The system's energy is much smaller than the bath's

Exponentiating both sides yields

$$\Omega_B(E_T - E) = \Omega_B(E_T) e^{-\beta E} \propto e^{-\beta E}$$

The system-dependent part is of greatest interest to us

$$P(\nu) \propto \Omega_B(E_T - E(\nu)) \propto e^{-\beta E}$$

(Marginalization) (Taylor expansion)

$P(\nu) \propto e^{-\beta E}$

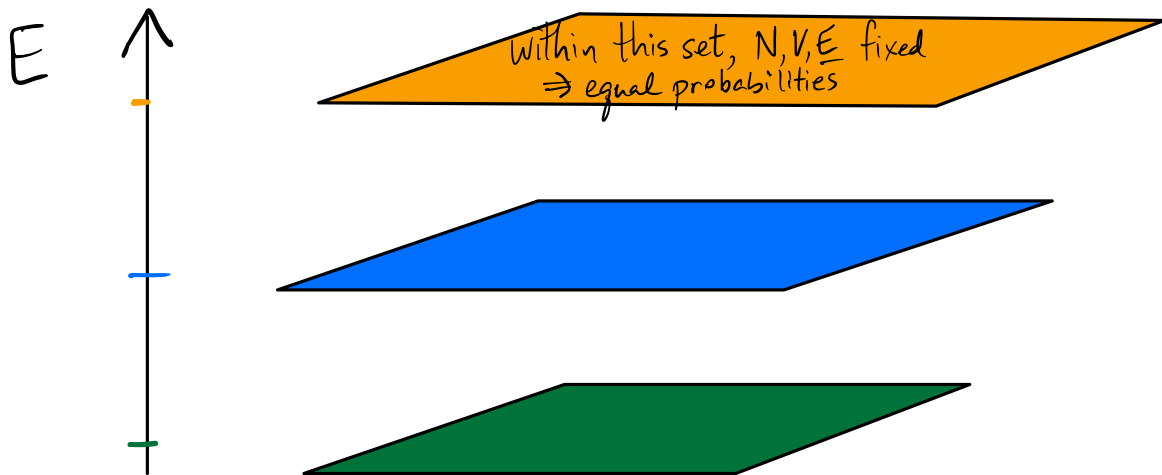
Boltzmann
Distribution!

E is no longer fixed. It fluctuates!

Are the system microstates ν equally likely?

In the microcanonical ensemble — constrained (N, V, E) — two microstates ν_1 and ν_2 had equal probability.

(ν_1 and ν_2 had to have equal energy due to the constraint.)



how many?
 $\Omega(N, V, E)$

Q: Why the $\propto$ sign? ($P(\nu) \propto e^{-\beta E}$) $P(\nu) = \frac{e^{-\beta E}}{Q}$

A: We would need to normalize the distribution
(Normalization Factor)

$$e^{-\beta E(\nu)}$$

Boltzmann Factor

Depends on β ? on ν ?

$$1 = \sum_{\nu} P(\nu) = \sum_{\nu} \frac{e^{-\beta E}}{Q} = \frac{1}{Q} \sum_{\nu} e^{-\beta E} \Rightarrow$$

$$Q(\beta) = \sum_{\nu} e^{-\beta E} \quad \text{dummy variable}$$

Depends on ν ?

No, ν is a dummy variable being summed over.

$$\Rightarrow P(\nu) = \begin{cases} \frac{e^{-\beta E(\nu)}}{Q(\beta, V, N)}, & \text{if } N_{\nu} = N, V_{\nu} = V \\ 0, & \text{otherwise} \end{cases}$$

What is $P(E)$?

All of those possible ν 's with $E_\nu = E$ have the same probability and combine together to give the probability of E .

$$P(E) = \sum_{\substack{\nu \text{ with} \\ E_\nu = E}} P(\nu) \propto e^{-\beta E} \sum_{\substack{\nu \text{ with} \\ E_\nu = E}} (1)$$

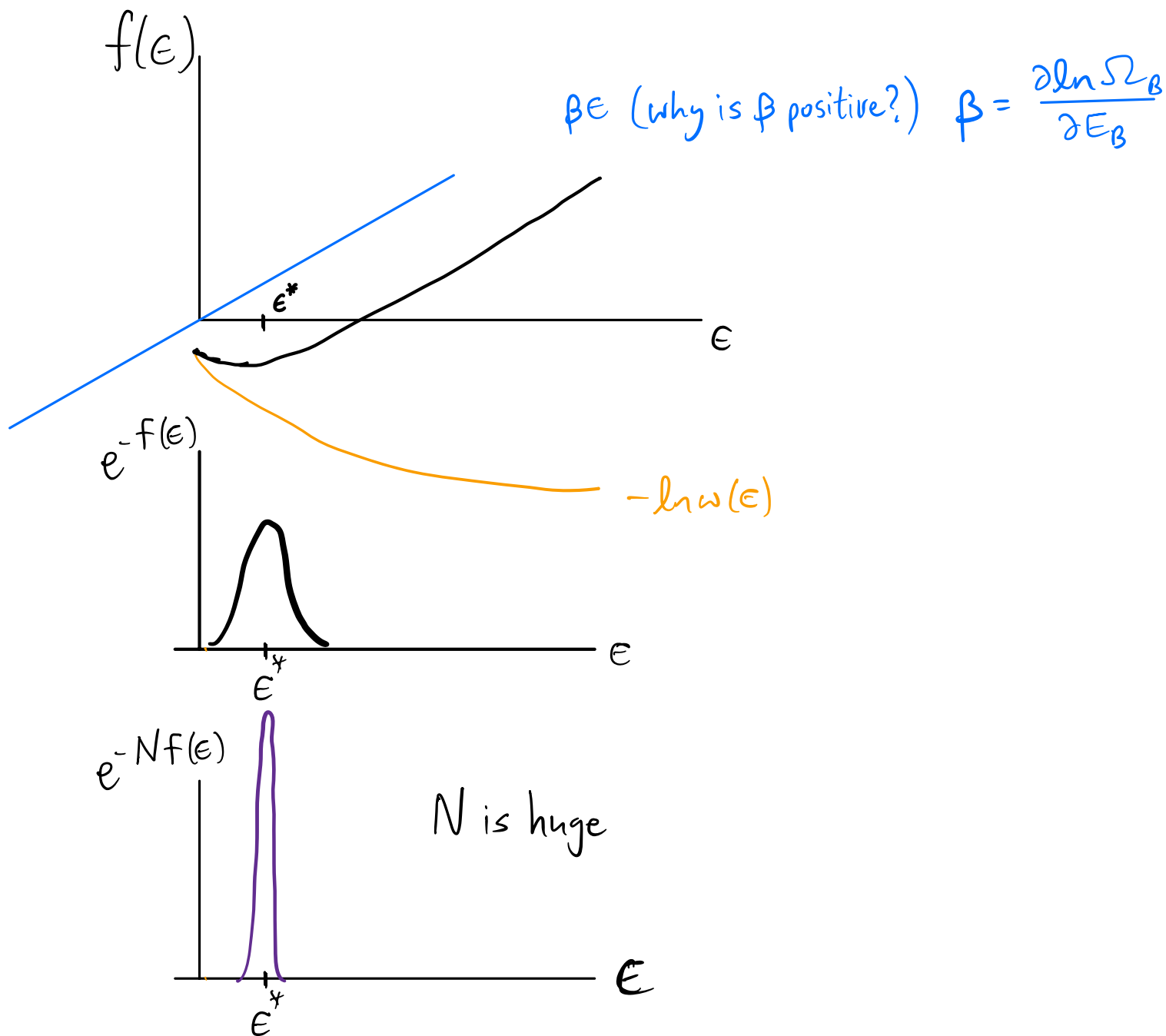
$$= e^{-\beta E} \Omega(E) e^{+N \ln w} \quad \text{"large deviation form"}$$

$$\Rightarrow P(E) \propto \exp \left[-N \left(\underbrace{\beta E}_{\text{extensive}} - \underbrace{\ln w}_{\text{intensive}} \right) \right] \quad \begin{array}{l} \text{another} \\ \text{"large deviation form"} \end{array}$$

$\left(\epsilon = \frac{E}{N} \right)$

Let's study this distribution in the limit of a big system.
($N \rightarrow \infty$)

$$\text{Define } f(\epsilon) = \beta \epsilon - \ln w \quad \Rightarrow P(E) \propto e^{-N f(\epsilon)}$$



In the macroscopic limit, only one value of $E = E^*(\beta)$ has non-negligible weight.

Things look boring at the macroscopic scale.

(no detected fluctuations)

We've said a little about the distribution $P(E)$.
How do we compute the typical E , i.e. $\langle E \rangle$?