

Lecture 12

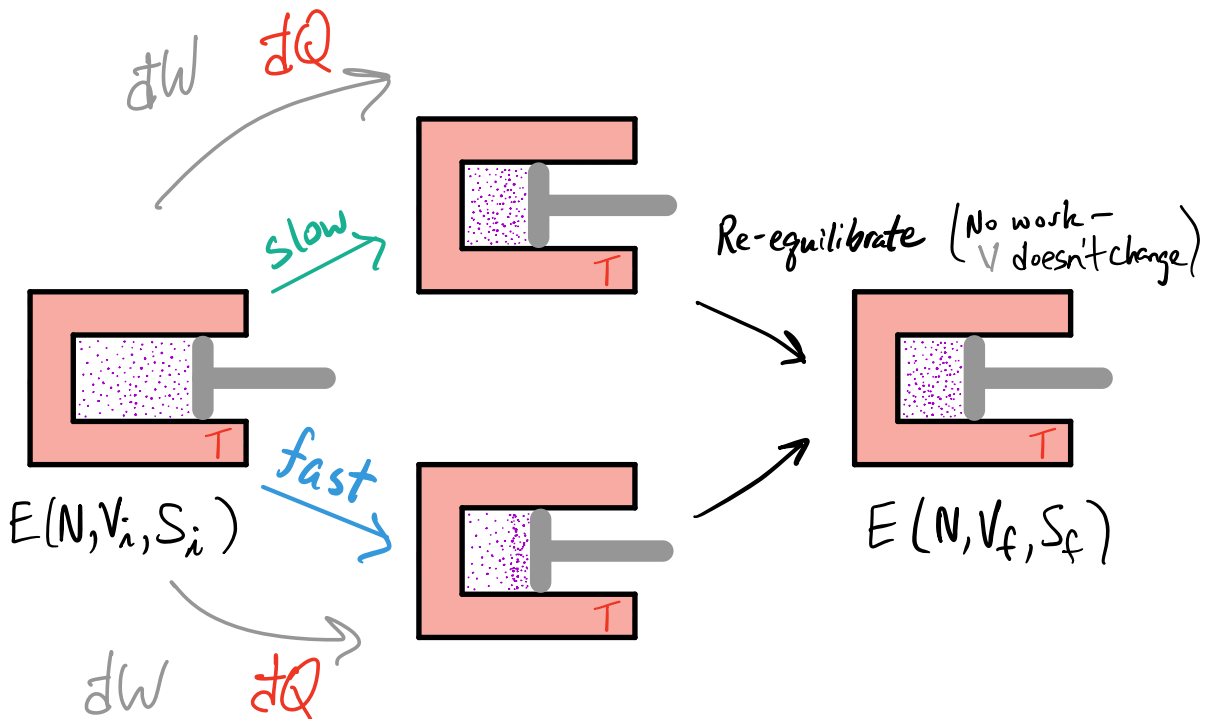
Recall from last lecture...

Ⓘ What is free about a free energy?

$$dE = \underline{TdS} - \underline{pdV} + \mu dN \rightarrow 0$$

Work: Change in energy due to variables I manipulate

Heat: Change in energy due to flows into or out of the thermal bath. We don't get to use these; they're along for the ride.



$$\int_{\text{fast}} dt \frac{dE(t)}{dt} = E(N, V_f, S_f) - E(N, V_i, S_i)$$

$$\int_{\text{slow}} dt \frac{dE(t)}{dt} = E(N, V_f, S_f) - E(N, V_i, S_i)$$

shorthand \Rightarrow

$$\int_{\text{fast}} dE = \int_{\text{slow}} dE$$

$$\Rightarrow \int_{\text{fast}} (\delta Q + \delta W) = \int_{\text{slow}} (\delta Q + \delta W)$$

(Sign convention: Positive W is work done on the system)
 Positive Q is heat flow into the system)

The fast process effectively pushes against more molecules

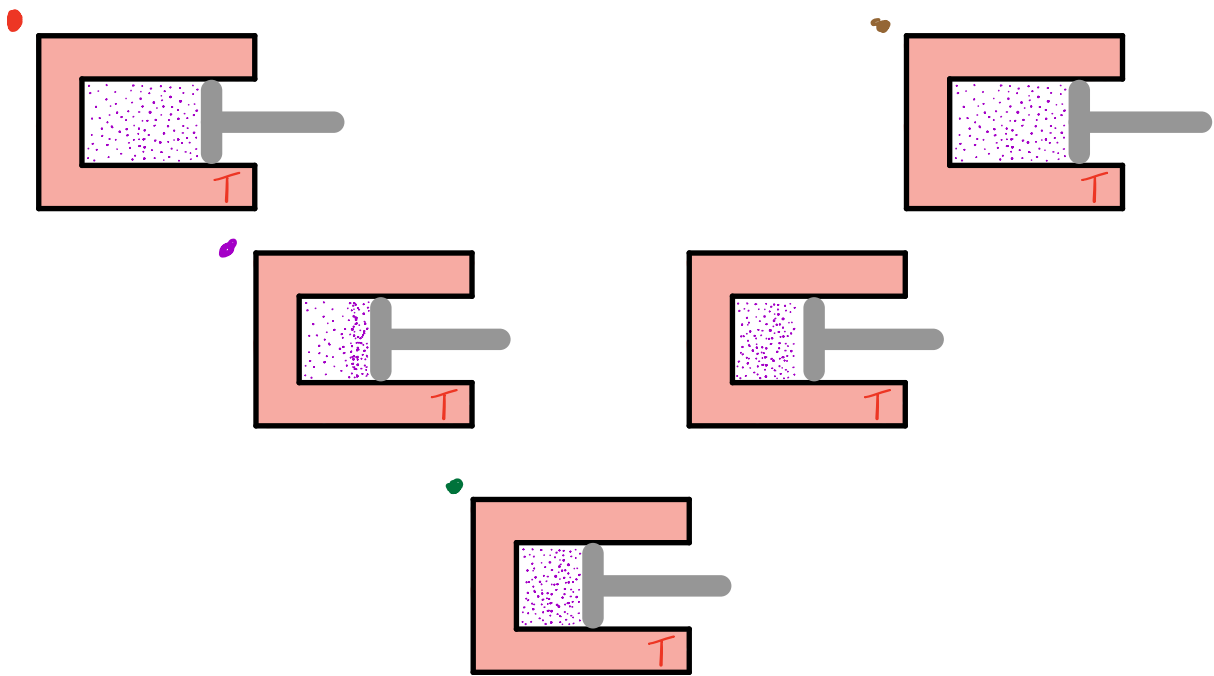
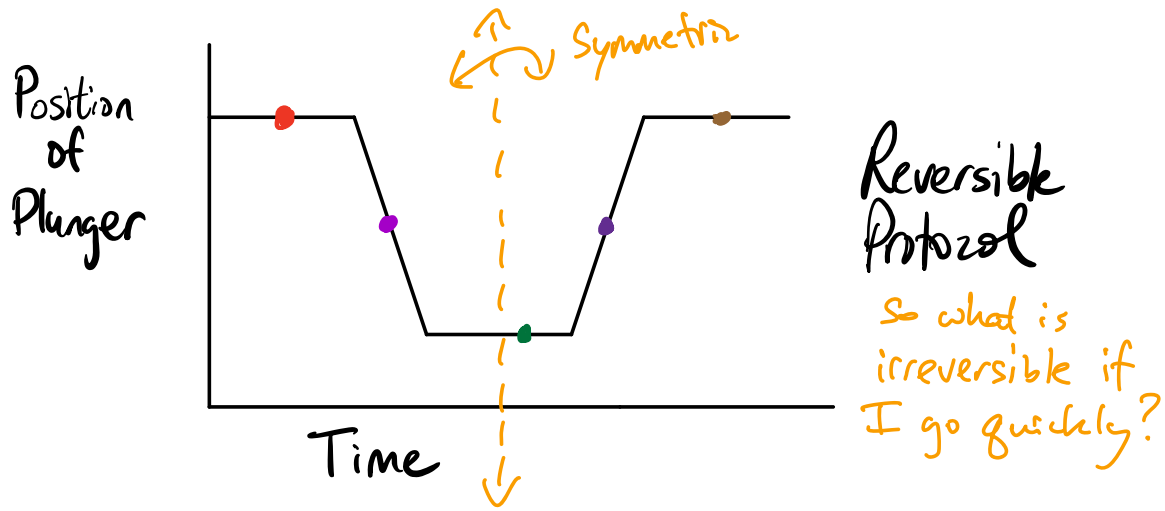
$$\Rightarrow \int_{\text{fast}} \delta W > \int_{\text{slow}} \delta W$$

Hence $\int_{\text{fast}} \delta Q < \int_{\text{slow}} \delta Q$
(δQ is negative)

Fast process has MORE heat flow out of the system.

The best we'll ever do is to go slowly (reversibly) so that the gas gets to re-equilibrate before I push any more.

Reversible vs Irreversible...



I visit different microstates unless I go slowly enough to visit equilibrium microstates in each direction

$$\boxed{dW \geq dW_{rev}}$$

A reversible transformation requires the least amount of work at fixed T .

If we go reversibly...

$$dW_{rev} = -pdV, \text{ so } dE = TdS + dW_{rev}$$

$$\Rightarrow dW_{rev} = dE - TdS$$

$$= d(E - TS) \Big|_T \leftarrow \text{Fixed } T$$

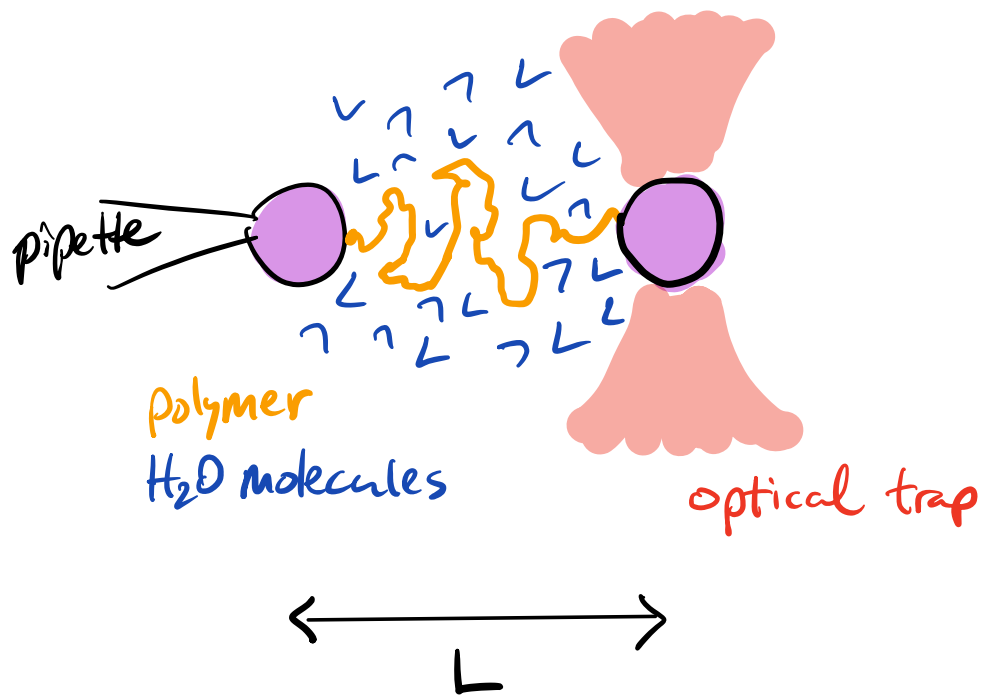
$$= (dA) \Big|_T \leftarrow \text{provided I am in contact w/ a fixed } T \text{ temp. bath}$$

$$\boxed{dW \geq (dA) \Big|_T}$$

The most work that can be extracted comes from a reversible transformation, in which case

$$W = \Delta A = A(N, V_f, T) - A(N, V_i, T)$$

These ideas extend beyond gasses and pistons.
They also apply to microscopic systems...



What am I going to measure? L

What am I not going to measure? H_2O molecules (bath)
interior monomers of the polymer

My microstates ν are given by the polymer configurations

$$\nu = (L, \{i\})$$

\uparrow end-to-end distance \uparrow everything else

$$P(v) = \frac{e^{-\beta E(v)}}{Q} \quad (\beta \text{ from } H_2O \text{ molecules})$$

What is $P(L)$? (Free polymer at this point - free to fluctuate that is)

$$P(L) \propto \sum_{\{i\}} e^{-\beta E(L, \{i\})}$$

Marginalize (average) over degrees of freedom I'm not measuring

$$\Rightarrow \ln P(L) = \ln(\text{const.}) + \ln \left(\sum_{\{i\}} e^{-\beta E(L, \{i\})} \right)$$

$$\Rightarrow \frac{\partial \ln P(L)}{\partial L} = \frac{\beta \sum_{\{i\}} \left(\frac{-dE(L, \{i\})}{dL} \right) e^{-\beta E(L, \{i\})}}{\sum_{\{i\}} e^{-\beta E(L, \{i\})}}$$

The statistical weight for microstate $\{i\}$ when L is fixed

$$P(\{i\} | \text{fixed } L) = \frac{e^{-\beta E(L, \{i\})}}{\sum_{\{i\}} e^{-\beta E(L, \{i\})}$$

Therefore

$$\frac{d \ln P(L)}{dL} = \beta \left\langle \frac{-dE(L, \xi_i)}{dL} \right\rangle_{\xi_i}$$

Averaging
over all possible
polymers with
length L

$$= \beta \left(\text{mean force by polymer} \right. \\ \left. \text{on coordinate } L \right)$$

$$= -\beta \left(\text{mean force applied to} \right. \\ \left. \text{the polymer to fix } L \right)$$

$$\int_{L_i}^{L_f} dL \frac{d \ln P(L)}{dL} = \beta \int_{L_i}^{L_f} \underbrace{dL}_{\text{distance}} \left\langle \underbrace{\frac{-dE(L, \xi_i)}{dL}}_{\text{mean force}} \right\rangle_{\xi_i}$$

$$\ln P(L_f) - \ln P(L_i)$$

$$\Rightarrow \ln P(L_f) - \ln P(L_i) = -\beta W_{\text{rev}}$$

Reversible Work Theorem:

$$\frac{P(L_f)}{P(L_i)} = e^{-\beta W_{\text{rev}}(L_i \rightarrow L_f)}$$
$$= e^{-\beta [A(L_f) - A(L_i)]}$$

Remember $\ln Q = -\beta A$, so we could also have seen this as

$$\frac{P(L_f)}{P(L_i)} = \frac{Q(L_f)}{Q(L_i)} = \frac{e^{-\beta A(L_f)}}{e^{-\beta A(L_i)}} = e^{-\beta W_{\text{rev}}(L_i \rightarrow L_f)}$$

(
 ratio of
 partition functions